Introduction to Thermoelectrics



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Thermoelectric Effect

Thermoelectric Effect: direct conversion of temperature differences to electric voltage and vice versa (Seebeck effect, Peltier effect, Thomson effect)

The Peltier effect

The Seebeck effect





Thermoelectric Coefficients



1) Seebeck coefficient (Thermopower, thermoelectric power, thermal EMF coefficient):

$$S_{AB} = \frac{V}{\Delta T}$$

2) Peltier coefficient (π): how much heat is carried per unit charge

$$\pi_{AB} = \frac{q}{I}$$
 q: rate of heating or cooling at each junction

Kelvin (Thomson) relation:

$$\pi_{AB}=S_{AB}T$$

3) Thomson coefficient (τ): a spatial gradient in temperature can result in a gradient in the Seebeck coefficient

$$\tau_A - \tau_B = T \frac{dS_{AB}}{dT}$$

Brief History of Thermoelectrics

Seebeck Effect (1821-3)





Thomas Johann Seebeck

Peltier Effect (1834)

In 1834, Jean Charles Athanase Peltier found that an electrical current would produce heating or cooling at the junction of two dissimilar metals.



In 1851 Gustav Magnus discovered the Seebeck voltage does not depend on the distribution of temperature along the metals between the junctions an indication that the thermopower is a thermodynamic state function.



Thomson Effect (1851)

In 1851, William Thomson (later Lord Kelvin) issued a comprehensive explanation of the Seebeck and Peltier Effects and described their interrelationship (known as the Kelvin Relations).

Gustav Magnus Basis for thermocouple

http://www.thermoelectrics.caltech.edu/thermoelectrics/history.html



Brief History of Thermoelectrics

Figure of Merit, ZT (1949)





In 1949 Abram Fedorovich loffe developed the modern theory of thermoelectricity using the concept of the 'figure of merit' zT.

One of the first demonstrations of 0 C cooling was by H. Julian Goldsmid in 1954 using thermoelements based on Bi_2Te_3 , and one of the first to utilize the thermoelectric quality factor, identifying the importance of high mobility and effective mass combination and low lattice thermal conductivity in semiconductors



In 1995, Glen Slack summarized the material requirements succinctly in the "phonon-glass electroncrystal" concept.

http://www.thermoelectrics.caltech.edu/thermoelectrics/history.html

Thermoelectric Refrigeration

Electric Voltage (Current) → Temperature Difference

Electric Current, *I* :

$$I=\frac{\sigma VA}{L}$$

Rate of Heat Flow, q:

$$q = -\frac{\kappa A \Delta I}{L}$$

 $\boldsymbol{L} \boldsymbol{\Lambda} \wedge \boldsymbol{T}$

Cooling Power, *q*₁ :

$$q_1 = \big(S_p - S_n\big)IT_1 - (T_2 - T_1)\big(K_p + K_n\big) - I^2(R_p + R_n)/2$$

Coefficient of Performance, COP, $\boldsymbol{\Phi}$:

$$\phi = \frac{\left(S_p - S_n\right)IT_1 - (T_2 - T_1)\left(K_p + K_n\right) - I^2(R_p + R_n)/2}{\left(S_p - S_n\right)I(T_2 - T_1) + I^2(R_p + R_n)}$$

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Current for the maximum cooling power, I_q :

$$I_q = (S_p - S_n)T_1/(R_p + R_n) \implies \phi_q = \frac{ZT_1^2/2 - (T_2 - T_1)}{ZT_2T_1}$$

For maximum COP,

$$I_{\phi} = \frac{(S_p - S_n)(T_2 - T_1)}{(R_p + R_n)\{(1 + ZT_m)^{1/2} - 1\}} \quad \phi_{max} = \frac{T_1 \left\{ (1 + ZT_m)^{1/2} - (\frac{T_2}{T_1}) \right\}}{(T_2 + T_1)\{(1 + ZT_m)^{1/2} + 1\}}$$

T)

Thermoelectric Refrigeration

Electric Voltage (Current) → Temperature Difference

Figure of Merit, *Z* :

$$Z = \frac{(S_p - S_n)^2}{(K_p + K_n)(R_p + R_n)} \longrightarrow \text{This should be as small as possible.}$$

So, the form factors satisfy the relation

$$\frac{L_n A_p}{L_p A_n} = \left(\frac{\rho_p k_n}{\rho_n k_p}\right)^{1/2}$$

$$Z = \frac{(S_p - S_n)^2}{\left\{ (k_p \rho_p)^{1/2} + (k_n \rho_n)^{1/2} \right\}^2}$$

The maximum temperature depression, ΔT_{max} , is reached when the cooling power and, thus, the COP fall to zero.

$$\Delta T_{max} = \frac{1}{2}ZT_1^2$$



Thermoelectric Refrigeration

Electric Voltage (Current) → Temperature Difference









$$\phi_{max} = \frac{T_1 \left\{ (1 + ZT_m)^{1/2} - (\frac{T_2}{T_1}) \right\}}{(T_2 + T_1) \{ (1 + ZT_m)^{1/2} + 1 \}}$$

Julian H. Goldsmid, Introduction to Thermoelectricity, Springer (2010)

Thermoelectric Generation

Temperature Difference → Electric Voltage (Current)

Thermal EMF,
$$E_{emf}$$
: $E_{emf} = (S_p - S_n)(T_h - T_c)$
Electric Current, I :
 $I = \frac{(S_p - S_n)(T_h - T_c)}{R_p + R_n + R_L}$
Power delivered to the load, w :
 $w = I^2 R_L = \left(\frac{(S_p - S_n)(T_h - T_c)}{R_p + R_n + R_L}\right)^2 R_L$
Total rate of heat flow from the source, q_I :
 $q_1 = (S_p - S_n)IT_1 + (K_p + K_n)(T_h - T_c)$

Efficiency, *η*:

$$\eta = w/q_1$$

*An increase in the load resistance reduces the power output but increases the efficiency.

For the maximum efficiency:

$$M = \frac{R_L}{R_p + R_n} = (1 + ZT_m)^{1/2}$$

Maximum Efficiency, *η*:

$$\eta = \frac{(T_h - T_c)(M - 1)}{T_h \left(M + \frac{T_c}{T_h}\right)}$$

*If ZT_m were much greater than unity, M would also be very large and the efficiency would approach $(T_h-T_c)/T_h$, which is the value for the **Carnot cycle**.

Thermoelectric Generation

Temperature Difference → Electric Voltage (Current)



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Thermoelectric Applications

Thermoelectric : simple structure and high reliable



TEC Applications



TEG Applications Using Waste Heat



- BMW, GM, Volkswagen

- 10% Fuel Efficiency Increase





Industrial TEG - Incineration Plant - Industrial facilities



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TEG Applications Using Waste Heat





TEG Applications



Portable TEG for outdoor



The maximum efficiency of a thermoelectric material for both power generation and cooling is determined by its figure of merit (zT):

Device:
$$\eta_{\text{max}} = \frac{\Delta T}{T_h} \cdot \frac{\sqrt{1+zT}-1}{\sqrt{1+zT}+T_c/T_h}$$
 $\frac{\Delta T}{T_h}$: Carnot efficiency







Breaking

the law

Coupling

Pisarenko Relation

$$S = \frac{8\pi^2 k_B T}{3qh^2} \cdot m_d^* \left(\frac{\pi}{3n}\right)^{2/3}$$

Wiedemann-Franz Law

$$\frac{k_e}{\sigma} = LT$$

Decoupling

PGEC (Phonon-Glass Electron-Crystal) Quantum Confinement Carrier Filtering Hierachical Structuring (defects) PGEC without cages

Phonon Scattering



Higher Seebeck coefficient (S) Higher electrical conductivity (σ) Lower thermal conductivity (k)

Engineering Decoupling

Glasses have low lattice thermal conductivity but low Seebeck coefficient. **Crystals** have high electrical conductivity but high thermal conductivity.

" Phonon-glass electron-crystal (PGEC)"

G. A. Slack, in CRC Handbook of Thermoelectrics (ed. Rowe, M.) 407-440 (CRC, Boca Ralon, 1995)

Original Paper: A. F. loffe, S. V. Airapetiants, A. F. loffe, N. V. Kolomoets and L. S. Stil'bans, *Dokl. Akad. Nauk. SSSR* 106,981 (1956).



$zT = \frac{S^2 \sigma T}{\sigma}$ • Electrical Conductivity Seebeck Coefficient $\sigma = en\mu$ Mott relation for degenerate statistics, Semiconductor -**Mott-transition** → Metal $S = \frac{\pi^2}{3} \cdot \frac{k_B}{e} \cdot k_B T \cdot \left| \frac{1}{g(E)} \cdot \frac{\partial g(E)}{\partial E} + \frac{1}{\mu(E)} \cdot \frac{\partial \mu(E)}{\partial E} \right|_{E=E_{-}}$ 1.5 k -80 $S = \frac{8\pi^2 k_B T}{3ah^2} \cdot m_d^* \left(\frac{\pi}{3n}\right)^{2/3}$ zT 1.0 Ω⁻¹cm⁻ Z 0.5 1000 B Na or Tl-doped PbTe 2% -2 1.5 g(E) 0.0 10¹⁹ 10²⁰ 10^{21} 10¹⁷ 10¹⁸ n_{c} (cm⁻³) 1.0 zT $-E_F \int E_R$ S. Lee et al., JECS 32 (2012) Ná Na 0.5 - Doping doped doped - Nonstoichiometry F 0.0 - Reduction 400 500 600 700 800 300 - Crystal Anisotropy Temperature (K) - Etc.

Heremans et al. Science 321, 554 (2008)





Phonon Scattering

- Complex Structure Multiple Sites
- Large Unit Volume
- Heavy Atoms
- Random Vacancies
- Local Lattice Distortion
- Low Symmetry
- Complex Compositions (Alloying, Solid Solutions)
- Rattling lons
- Segregation of Chemistry-Second Phases
- Nanostructures
- Intrinsic Localized Modes (ILM)
- Etc.



 $S^2 \sigma T$





Chemical Exfoliation of Bi₂Te₃

- Lithium Intercalation: $Bi_2Te_3 + x$ (n-BuLi) $\Rightarrow Li_xBi_2Te_3 + x/2$ octane
- \bigcirc Exfoliation : $Li_xBi_2Te_3 + xH_2O \Rightarrow Bi_2Te_3$ (exfoliated) + xLiOH + x/2H₂ \uparrow



J. Y. Kim et al. J. Mater. Chem. A 2013

Thermoelectric Materials





Weak coupling of "Rattling filler" with the SKD.

Decoupling: High Electrical Conductivity and low Thermal Conductivity



Filled Skutterudites



Xun Shi, et al., J. Am. Chem. Soc. 2011, 133, 7837–7846



Broad Spectrum of Localized Modes

TABLE II. Spring constant k and resonance frequency ω_0 in the [111] and [100] directions of $R_{0.125}Co_4Sb_{12}$, where R = Ia, Ce, Eu, Yb, Ba, Sr, Na, and K.

		[111]		[100]	
R	Mass (10 ⁻²⁶ Kg)	k (N/m)	$\omega_0 (\text{cm}^{-1})$	k (N/m)	$\omega_0 (\text{cm}^{-1})$
La	23.07	36.10	66	37.42	68
Се	23.27	23.72	54	25.18	55
Eu	25.34	30.16	58	31.37	59
Yb	28.74	18.04	42	18.88	43
Ва	22.81	69.60	93	70.85	94
Sr	14.55	41.62	90	42.56	91
Na	3.819	16.87	112	17.18	113
К	6.495	46.04	141	46.70	142

1. J. Yang, W. Zhang, S. Q. Bai, Z. Mei, and L. Chen, Appl. Phys. Lett. 90, 192111 (2007)

Electronegativity Rule for SKD Filler

i: K, Ba, Sr, La, Yb, Na, Ce, etc. **Provide and Provide And Provid**

j: Br, Cl, Se, S, etc.

thermal conductivity

*The *j* guests show strong covalent bonding, leading to cluster vibration which decreases

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Ionic Size Rule for SKD Filler

Ionic Radii vs Atomic Radii

Half-Heusler

 $(Hf_{0.5}Zr_{0.5})_{1-x}Ti_{x}NiSn_{0.998}Sb_{0.002}$

Charge Transport in Bi₂Te₃ System

- Bi_2Te_3 is a unique thermoelectric material which can be used around room temperature.
- It can be used for both cooling and power generation applications.
- But it has weak Van der Waals bonding. → need powder processing

For non-degenerate semiconductor thermoelectrics

$$ZT = \frac{[\eta - (r + 5/2)]^2}{(\beta \exp(\eta))^{-1} + (r + 5/2)}$$

η: Fermi energyr: scattering parameterβ: materials parameter

 $\boldsymbol{\beta}$ was first introduced by Chasmar and Stratton,

$$\beta = \left(\frac{k_B}{e}\right)^2 \frac{\sigma_0 T}{k_L} \qquad \sigma_0 = 2e\mu \left(\frac{2\pi m^* k_B T}{h^2}\right)^{3/2}$$
$$\beta \propto \left(\frac{\mu}{k_L}\right) \left(\frac{m^*}{m}\right)^{3/2}$$

*Although **ZT** equation holds only for a non-degenerate semiconductor, the materials parameter, β , remains useful when the material is partly or completely degenerated.

Julian H. Goldsmid, *Introduction to Thermoelectricity*, Springer (2010)

Oxide Thermoelectric Materials

M. Ohtaki, Global COE Program Novel Carbon Resources Sciences Newsletter, 2010. 05.

N. V. Nong et al., Adv. Mater. 23, 2484 (2011)

Oxide Thermoelectric Materials

How to create charge carriers in oxides?

S. Lee et al., JECS 32 (2012)

Oxide Thermoelectric Materials

How to generate the defects?

Equilibrium Concentration of Point Defects

Equilibrium Concentration varies with temperature and atmosphere:

Oxide Thermoelectric Materials

How to generate the charge carriers?

For *n*-type: Donor-doped MO / Reduced MO

Oxide Thermoelectric Materials

Oxide Thermoelectric Materials

Quantum Confinement Effect

Conductive AFM

Textured Ceramic

Randomly Oriented SBN [001] Textured SBN B&C Power Factor, S²σ (μW/cmK²) - A: Randomly oriented SBN ceramic - B: Textured SBN ceramic (measured $\perp c$ -axis growth) C: Textured SBN ceramic (measured //c-axis growth) 450 500 150 200 250 300 350 400 550 **Temperature (K)**

S. Lee et al., JMR 26[1], 26-30 (2011)

Thermoelectric Modules

Thermoelectric Module

Engineering Thermoelectric Modules

Engineering Thermoelectric Modules

Engineering Thermoelectric Modules

Bottleneck of TEG Module for Mid-High Temperature Applications

Development of Thermoelectric Modules

Diffusion Barrier Material and Process Design

- Development of Diffusion Barrier (DB) Materials
- Development of DB Process

DB by TS

Development of Thermoelectric Modules

High Temperature Stability

- Enhancing Technology of high temperature stability (Oxidation, Evaporation)
- Module Packaging Technology

Conclusion

Thank you for your attention!

